# Painlevé Analysis and Some Solutions of $(2+1)$-Dimensional Generalized Burgers Equations 

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#### Abstract

Burgers equation $u_{t}=2 u u_{x}+u_{x x}$ describes a lot of phenomena in physics fields, and it has attracted much attention. In this paper, the Burgers equation is generalized to $(2+1)$ dimensions. By means of the Painlevé analysis, the most generalized Painlevé integrable (2+1)-dimensional integrable Burgers systems are obtained. Some exact solutions of the generalized Burgers system are obtained via variable separation approach.


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Key words: Burgers equation, variable separation

## 1 Introduction

The Burgers equation $u_{t}=2 u u_{x}+u_{x x}$ has attracted much attention since it was first proposed by Bateman. ${ }^{[1]}$ Then Burgers ${ }^{[2]}$ gave its some special solutions in 1940. Later on, Cole and Hopf ${ }^{[3,4]}$ independently pointed out that any of the solutions of the heat equation $\xi_{t}=\xi_{x x}$ can be mapped to a solution of the Burgers equation. Recently, the Painlev́e property (PP) has shown its usefulness in studying of partial differential equation (PDE), especially in the sense of integrability. Some connections of PP with inverse scattering transform (IST) were found. Through Painlev́e analysis for an integrable model, we could probably find its various interesting properties like the Lax pairs. ${ }^{[5-7]}$ In this paper, we would like to find $(2+1)$-dimensional Burgers equations which are Painlev́e integrable, and then use the variable separation approach to get some of their solutions.

## 2 Generalized Burgers Equations

The most generalized Burgers equation may have the form

$$
\begin{align*}
\frac{\partial}{\partial t} u= & a_{1} u \frac{\partial}{\partial x} u+a_{2} \frac{\partial^{2}}{\partial x^{2}} u+a_{3} u \frac{\partial}{\partial y} u+a_{4} \frac{\partial^{2}}{\partial y^{2}} u \\
& +a_{5} \frac{\partial}{\partial x} u+a_{6} \frac{\partial}{\partial y} u+a_{7} u+a_{8} \frac{\partial^{2}}{\partial x \partial y} u \\
& +b_{1} v \frac{\partial}{\partial x} v+b_{2} \frac{\partial^{2}}{\partial x^{2}} v+b_{3} v \frac{\partial}{\partial y} v(x, y, t) \\
& +b_{4} \frac{\partial^{2}}{\partial y^{2}} v+b_{5} \frac{\partial}{\partial x} v+b_{6} \frac{\partial}{\partial y} v+b_{7} v \\
& +b_{8} \frac{\partial^{2}}{\partial x \partial y} v+c_{1} u v_{y}+c_{2} v u_{x}+c_{3} u v  \tag{1}\\
\frac{\partial}{\partial y} u= & \frac{\partial}{\partial x} v \tag{2}
\end{align*}
$$

where $a_{i}, b_{i}$, and $c_{i}$ are all constants, which means that we are finding constant coefficient integrable models of Burger's type.

To check whether the given equations have Painlev́e property, we have several choices, such as ARS (Ablowitz-Ramani-Segur) algorithm, ${ }^{[8]}$ the WTC (Weiss-TaborCarnevale) approach, ${ }^{[9]}$ Kruskal's simplification, ${ }^{[10]}$ the Conte's invariant method, ${ }^{[5]}$ Pickering's approach, ${ }^{[7]}$ and

Lou's extended method. ${ }^{[11]}$ Here we used WTC approach to analyse the equations.

As usual, we use a simple three-step version to find the possible models which have PP.
(i) Leading order analysis. Letting

$$
\begin{align*}
& u=\phi^{\alpha} \sum_{j=0}^{\infty} u_{j} \phi^{j}  \tag{3}\\
& v=\phi^{\beta} \sum_{j=0}^{\infty} v_{j} \phi^{j} \tag{4}
\end{align*}
$$

and substituting them into Eqs. (1) and (2) and balancing the leading nonlinear and leading dispersive terms, we could easily get $\alpha=\beta=-1$.
(ii) To find the resonance points. Because we are treating them as the Burger's type equation, we suppose that both $a_{1}$ and $a_{2}$ are not equal to zero. And the generalized ( $2+1$ )-dimensional Burgers equations have the resonance points which the original Burgers equation $u_{t}=2 u u_{x}+u_{x x}$ has. That means the generalized Burgers equations must have the resonance points at $j=2$. Using the computer algebra such as Maple, we found that there are several cases in which the generalized $(2+1)$ dimensional Burgers equations are Painlev́e integrable. And the resonance points are $j=2, j=1$ and $j=-1$. Because of the complexity of the results, we do not write them down here.
(iii) To check if the compatibility at the resonance points holds, by which we mean that $u_{1}, u_{2}$ and $\phi$ must be arbitrary.

After finishing the above three steps and using rescaling procedure, we get the final nontrivial generalized $(2+1)$-dimensional Painlev́e integrable Burgers equations

$$
\begin{equation*}
u_{t}=a_{1} u u_{x}+a_{2} u_{x x}+b_{1} v v_{x}+\frac{a_{2} b_{1}}{a_{1}} v_{x y}, \quad u_{y}=v_{x} . \tag{5}
\end{equation*}
$$

## 3 Some Solutions of Generalized Burgers Equations

Now we try to find some solutions of Eq. (5) through the variable separation approach. ${ }^{[11]}$ We rewrite Eq. (5) as

$$
u_{t}=A_{1} u u_{x}+A_{2} u_{x x}+B_{1} v v_{x}+\frac{A_{2} B_{1}}{A_{1}} v_{x y}
$$

[^0]\[

$$
\begin{equation*}
u_{y}=v_{x} \tag{6}
\end{equation*}
$$

\]

The first step of the approach is to suppose that $u$ and $v$ take the following forms
$u=2\left(A_{2} / A_{1}\right) \ln (f)_{x}+u_{0}, \quad v=2\left(A_{2} / A_{1}\right) \ln (f)_{y}+v_{0}, \quad(7)$ which can be found from the truncated Painlevé expansion, where $u_{0}, v_{0}$ are arbitrary known solutions of Eqs. (6). Here, however, we take $u_{0}=0$ and $v_{0}=v_{0}(y, t)$
for simplicity. The function $v_{0}$ is an arbitrary function of the indicated variables.

Then we suppose that $f$ has the form

$$
\begin{equation*}
f=1+a_{1} p(x, t)+a_{2} q(y, t)+A p(x, t) q(y, t) \tag{8}
\end{equation*}
$$

where $a_{1}, a_{2}$ and $A$ are constants to be determined; $p(x, t)$ and $q(y, t)$ are at present supposed to be arbitrary. Substituting Eqs. (7) with Eq. (8) into Eqs. (6) yields

$$
\begin{align*}
0= & 2 A_{2}\left(-A_{1} A^{2} p_{x} q^{2}-A_{1} a_{1}^{2} p_{x}-2 A_{1} a_{1} p_{x} A q\right) p_{t}+2 A_{2}\left(A_{1} A p_{x}-a_{1} A_{1} p_{x} a_{2}\right) q_{t} \\
& +2 A_{2}\left(2 a_{1} A_{1} A q p+A_{1} A q^{2} a_{2}+A_{1} a_{1}^{2} p+a_{1} A_{1}+a_{1} A_{1} a_{2} q+A_{1} A q+A_{1} A^{2} q^{2} p\right) p_{x t} \\
& +2 A_{2}\left(a_{2} A_{2} A_{1} A q^{2}+A_{2} A_{1} A^{2} q^{2} p+2 A_{2} a_{1} A_{1} A p q+a_{2} a_{1} A_{1}+a_{2} A_{2} a_{1} A_{1} q+A_{2} A_{1} A q+A_{2} A_{1} a_{1}^{2} p\right) p_{x x x} \\
& +2 A_{2}\left(-A_{2} A_{1} A^{2} p_{x} q^{2}-A_{2} A_{1} a_{1}^{2} p_{x}-2 A_{2} a_{1} A_{1} p_{x} A q\right) p_{x x} \\
& +2 A_{2}\left(-A_{2} a_{2} B_{1} a_{1} p_{x}+A_{2} B_{1} A p_{x}\right) q_{y y}+2 A_{2}\left(-B_{1} A_{1} a_{1} a_{2} p_{x} q_{y}+B_{1} A_{1} A p_{x} q_{y}\right) v_{0} . \tag{9}
\end{align*}
$$

Using the computer algebras such as Maple or Mathematica, we get

$$
\begin{align*}
0= & 2 J A_{1} A_{2}\left(-J p_{x}+f \partial_{x}\right)\left(p_{t}+A_{2} p_{x x}\right) \\
& -2 A_{2} K p_{x}\left(A_{2} B_{1} q_{y y}+A_{1} q_{t}+A_{1} B_{1} q_{y} v_{0}\right), \tag{10}
\end{align*}
$$

where $J=A q+a_{1}$ and $K=a_{1} a_{2}-A$. Because $p$ is $y$ independent and $q$ is $x$-independent, equation (10) can be separated into two parts,

$$
\begin{align*}
& p_{t}+A_{2} p_{x x}=d_{2} p^{2}+d_{1} p+d_{0}  \tag{11}\\
& A_{2} B_{1} q_{y y}+A_{1} q_{t}+A_{1} B_{1} q_{y} v_{0}=c_{2} q^{2}+c_{1} q+c_{0} \tag{12}
\end{align*}
$$

Substituting Eqs. (11) and (12) into Eq. (10), we can determine the values of $c_{i}$ and $d_{i}$. There are two cases. The first one is

$$
\begin{equation*}
d_{2}=0, \quad d_{0}=d_{1} / a_{1}, \quad A=a_{1} a_{2} \tag{13}
\end{equation*}
$$

with $a_{1}, a_{2}, c_{0}(y, t), c_{1}(y, t), c_{2}(y, t)$, and $d_{1}(t)$ being arbitrary constants/functions with indicated variables. The second one is

$$
\begin{align*}
& d_{2}=0, \quad c_{1}=\left(2 a_{1} c_{2}-A_{1} A d_{1}\right) / A \\
& d_{0}=\left(A_{1} A a_{2} d_{1}-a_{1} a_{2} c_{2}+A c_{2}\right) / A_{1} A^{2} \\
& c_{0}=a_{1}\left(a_{1} c_{2}-A_{1} A d_{1}\right) / A^{2} \tag{14}
\end{align*}
$$

with $A, a_{1}, a_{2}$, and $c_{2}(t)$ and $d_{1}(t)$ being arbitrary.
Both $v_{0}(y, t)$ and $q(y, t)$ are previously supposed to be arbitrary of the indicated variables. But from Eq. (12) we know that we must fix one of them, say $v_{0}$, and let the other be free. Hence, we obtain a solution of the generalized Burgers equation (6) with $q(y, t)$ being arbitrary

$$
u=\frac{2 A_{2}}{A_{1}} \frac{p_{x}\left(a_{1}+A q\right)}{\left(1+a_{1} p+a_{2} q+A p q\right)}
$$

$$
\begin{equation*}
v=\frac{2 A_{2}}{A_{1}} \frac{q_{y}\left(a_{2}+A p\right)}{\left(1+a_{1} p+a_{2} q+A p q\right)}+v_{0} \tag{15}
\end{equation*}
$$

where $p$ must satisfy Eq. (11) while $v_{0}$ is determined by $q(y, t)$ and $c_{i}(i=0,1,2)$ according to Eq. (12).

## 4 Summary and Discussion

In summary, using the WTC method, we obtained the most generalized $(2+1)$-dimensional constant coefficients Burgers equations which have Painlev́e property (Painlev́e integrable). Then with the help of Lou's separation method, ${ }^{[12]}$ we get solutions with some arbitrary functions. As in the articles, ${ }^{[12]}$ by fixing these arbitrary functions, we could get some interesting special solutions. We leave it undone because of two points. First, there are complete details in how to select functions to get interesting different solutions, which show the probable connections between non-integrable systems and integrable systems. Second, those who are at liberation know better how to fix the arbitrary functions to explain phenomena in experiments.

The next question is what type of partial differential equations can be separated in Lou's method to get solutions with arbitrary functions. This is investigated now by a group under professor S.Y. Lou. Another question is why those equations can be separated in this way. What is the mathematical or physical reasons? It deserves consideration.

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